# GENERALIZED NEHARI MANIFOLD AND SEMILINEAR SCHRÖDINGER EQUATION WITH WEAK MONOTONICITY CONDITION ON THE NONLINEAR TERM 

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#### Abstract

We study the Schrödinger equations $-\Delta u+V(x) u=f(x, u)$ in $\mathbb{R}^{N}$ and $-\Delta u-\lambda u=f(x, u)$ in a bounded domain $\Omega \subset \mathbb{R}^{N}$. We assume that $f$ is superlinear but of subcritical growth and $u \mapsto f(x, u) /|u|$ is nondecreasing. In $\mathbb{R}^{N}$ we also assume that $V$ and $f$ are periodic in $x_{1}, \ldots, x_{N}$. We show that these equations have a ground state and that there exist infinitely many solutions if $f$ is odd in $u$. Our results generalize those in [11] where $u \mapsto f(x, u) /|u|$ was assumed to be strictly increasing. This seemingly small change forces us to go beyond methods of smooth analysis.


## 1. INTRODUCTION

We consider the semilinear Schrödinger equations

$$
\begin{equation*}
-\Delta u+V(x) u=f(x, u), \quad u \in H^{1}\left(\mathbb{R}^{N}\right) \tag{1.1}
\end{equation*}
$$

and

$$
\begin{equation*}
-\Delta u-\lambda u=f(x, u), \quad u \in H_{0}^{1}(\Omega), \tag{1.2}
\end{equation*}
$$

where $\Omega \subset \mathbb{R}^{N}$ is a bounded domain and $H^{1}\left(\mathbb{R}^{N}\right), H_{0}^{1}(\Omega)$ are the usual Sobolev spaces. In both problems we make the following assumptions on $f$ :
$\left(F_{1}\right) f$ is continuous and $|f(x, u)| \leq C\left(1+|u|^{p-1}\right)$ for some $C>0$ and $p \in\left(2,2^{*}\right)$, where $2^{*}:=2 N /(N-2)$ if $N \geq 3$ and $2^{*}:=+\infty$ if $N=1$ or 2,
( $F_{2}$ ) $f(x, u)=o(u)$ uniformly in $x$ as $u \rightarrow 0$,
$\left(F_{3}\right) F(x, u) / u^{2} \rightarrow \infty$ uniformly in $x$ as $|u| \rightarrow \infty$, where $F(x, u):=\int_{0}^{u} f(x, s) d s$,
$\left(F_{4}\right) u \mapsto f(x, u) /|u|$ is non-decreasing on $(-\infty, 0)$ and on $(0, \infty)$.
The assumptions $\left(F_{1}\right)-\left(F_{3}\right)$ appear in [11] while a condition corresponding to $\left(F_{4}\right)$ is a little stronger there:
$\left(F_{4}^{\prime}\right) u \mapsto f(x, u) /|u|$ is strictly increasing on $(-\infty, 0)$ and on $(0, \infty)$.

[^0]As we shall see, this slightly weaker hypothesis will force us to go beyond methods of smooth analysis, and introducing a non-smooth approach in this context is in fact our main purpose. In what follows we shall frequently refer to different results and arguments in $[11,12]$. When such reference is made, it should be understood that no stronger conditions than $\left(F_{1}\right)-\left(F_{4}\right)$ were needed there.

The main results of this paper are the following two theorems:
Theorem 1.1. Suppose $f$ satisfies $\left(F_{1}\right)-\left(F_{4}\right), V$ and $f$ are 1-periodic in $x_{1}, \ldots, x_{N}$ and $0 \notin \sigma(-\Delta+V)$, where $\sigma(\cdot)$ denotes the spectrum in $L^{2}\left(\mathbb{R}^{N}\right)$. Then equation (1.1) has a ground state solution. If moreover $f$ is odd in $u$, then equation (1.1) has infinitely many pairs of geometrically distinct solutions.

Theorem 1.2. (i) Suppose $f$ satisfies $\left(F_{1}\right)-\left(F_{4}\right)$ and $\lambda \neq \lambda_{k}$ for any $k$, where $\lambda_{k}$ is the $k$-th eigenvalue of $-\Delta$ in $H_{0}^{1}(\Omega)$. Then equation (1.2) has a ground state solution. If moreover $f$ is odd in $u$, then equation (1.1) has infinitely many pairs of geometrically distinct solutions $\pm u_{k}$ such that the $L^{\infty}(\Omega)$-norm of $u_{k}$ tends to infinity with $k$.
(ii) If $\lambda=\lambda_{k}$ for some $k$, then the above results remain valid under the additional assumption that $f(x, u) \neq 0$ unless $u=0$.

Similar results, but under the stronger condition $\left(F_{4}^{\prime}\right)$, have been proved in [11].
As usual, a ground state is a solution which minimizes the functional corresponding to the problem over the set of all nontrivial $(u \neq 0)$ solutions. Later in this section we shall define what we mean by geometrically distinct solutions.

Existence of a ground state solution under the assumptions of Theorem 1.1 has been shown by S. Liu in [7]; since this result is an easy consequence of our approach, we include it here anyway. See also [16] where a number of results on ground states for problems similar to (1.1) and (1.2) has been proved and [13] where $\left(F_{4}\right)$ has been further weakened. Existence of ground states for systems of equations has been discussed in [8]. Concerning existence of infinitely many solutions we know of a result by Tang [14] where a condition different from $\left(F_{4}\right)$ has been introduced for (1.2), and by Zhong and Zou [16] where (1.1) and (1.2) have been considered under the same hypotheses as in Theorems 1.1 and 1.2. However, they needed an additional assumption which is not easy to verify unless $u \mapsto f(x, u) /|u|$ is "most times" strictly increasing.

Consider equation (1.1) under the assumptions of Theorem 1.1. Let $E:=H^{1}\left(\mathbb{R}^{N}\right)$. The functional corresponding to (1.1) is

$$
\Phi(u):=\frac{1}{2} \int_{\mathbb{R}^{N}}\left(|\nabla u|^{2}+V(x) u^{2}\right) d x-\int_{\mathbb{R}^{N}} F(x, u) d x
$$

It is well known (see e.g. [15]) that $\Phi \in C^{1}(E, \mathbb{R})$ and critical points of $\Phi$ are solutions for (1.1). Let $E=E^{+} \oplus E^{-}$be the decomposition corresponding to the positive and the negative part of the spectrum of $-\Delta+V$. Since $0 \notin \sigma(-\Delta+V)$, there exists an equivalent
inner product $\langle.,$.$\rangle in E$ such that

$$
\begin{equation*}
\Phi(u)=\frac{1}{2}\left\|u^{+}\right\|^{2}-\frac{1}{2}\left\|u^{-}\right\|^{2}-\int_{\mathbb{R}^{N}} F(x, u) d x \tag{1.3}
\end{equation*}
$$

where $u^{ \pm} \in E^{ \pm}$.
For equation (1.2) under the assumptions of Theorem 1.2 we put $E=H_{0}^{1}(\Omega)$ and we have the spectral decomposition $E=E^{+} \oplus E^{0} \oplus E^{-}$, where $E^{0}$ is the nullspace of $-\Delta-\lambda$ in $E$ and $0 \leq \operatorname{dim}\left(E^{0} \oplus E^{-}\right)<\infty$. Also here we can choose an equivalent inner product such that the corresponding functional $\Phi$ is of the form (1.3), with $\mathbb{R}^{N}$ replaced by $\Omega$.

The following set introduced by Pankov [9] is called the generalized Nehari manifold or the Nehari-Pankov manifold:

$$
\begin{equation*}
\mathcal{M}:=\left\{u \in E \backslash\left(E^{0} \oplus E^{-}\right): \Phi^{\prime}(u) u=0 \text { and } \Phi^{\prime}(u) v=0 \text { for all } v \in E^{0} \oplus E^{-}\right\} \tag{1.4}
\end{equation*}
$$

( $E^{0}$ is necessarily trivial in Theorem 1.1). ( $F_{4}$ ) implies $f(x, u) u \geq 0$, and the assumptions of Theorem 1.2 imply that if $\operatorname{dim} E^{0}>0$, then $f(x, u) u>0$ for $u \neq 0$. Hence $\mathcal{M}$ contains all nontrivial critical points of $\Phi$. Note that if $E^{0} \oplus E^{-}=\{0\}$, then $\mathcal{M}$ is the usual Nehari manifold [12]. Since this case is considerably easier to handle, we assume in what follows that $\sigma(-\Delta+V) \cap(-\infty, 0) \neq \emptyset$ in Theorem 1.1 and $\lambda \geq \lambda_{1}$ in Theorem 1.2. As in [11], for $u \notin E^{0} \oplus E^{-}$we define

$$
\begin{align*}
E(u):= & E^{0} \oplus E^{-} \oplus \mathbb{R} u=E^{0} \oplus E^{-} \oplus \mathbb{R} u^{+}  \tag{1.5}\\
& \text {and } \widehat{E}(u):=E^{0} \oplus E^{-} \oplus \mathbb{R}^{+} u=E^{0} \oplus E^{-} \oplus \mathbb{R}^{+} u^{+}
\end{align*}
$$

where $\mathbb{R}^{+}=[0, \infty)$. It has been shown there that if $\left(F_{4}\right)$ is replaced by $\left(F_{4}^{\prime}\right)$, then $\widehat{E}(u)$ intersects $\mathcal{M}$ at a unique point which is the unique global maximum of $\left.\Phi\right|_{\widehat{E}(u)}$. It has been shown in [16] by an explicit example that if $\left(F_{4}\right)$ but not $\left(F_{4}^{\prime}\right)$ holds, then (in the framework of Theorem 1.2) $\widehat{E}(u)$ and $\mathcal{M}$ may intersect on a finite line segment. In the next section we shall show that $\widehat{E}(u) \cap \mathcal{M} \neq \emptyset$ and if $w \in \widehat{E}(u) \cap \mathcal{M}$, then there exist $\sigma_{w}>0, \tau_{w} \geq \sigma_{w}$ such that $\widehat{E}(u) \cap \mathcal{M}=\left[\sigma_{w}, \tau_{w}\right] w$. In other words, $\widehat{E}(u) \cap \mathcal{M}$ is either a point or a finite line segment. We also show that a point $\widetilde{w} \in\left[\sigma_{w}, \tau_{w}\right] w$ is critical for $\Phi$ if and only if the whole segment $\left[\sigma_{w}, \tau_{w}\right] w$ consists of critical points.

In Theorem 1.1 the functional $\Phi$ is invariant with respect to the action of $\mathbb{Z}^{N}$ given by the translations $k \mapsto u(\cdot-k), k \in \mathbb{Z}^{N}$. Hence if $u$ is a solution of $(1.1)$, then so is $u(\cdot-k)$. This and the preceding paragraph justify the following definition: Two solutions $u_{1}$ and $u_{2}$ are called geometrically distinct if $u_{2} \neq u_{1}(\cdot-k)$ for any $k \in \mathbb{Z}^{N}$ and $u_{2} \notin\left[\sigma_{u_{1}}, \tau_{u_{1}}\right] u_{1}$. In Theorem 1.2 there is no $\mathbb{Z}^{N}$-invariance but we still want to identify solutions in $\widehat{E}(u) \cap \mathcal{M}$. So $u_{1}, u_{2}$ are geometrically distinct if $u_{2} \notin\left[\sigma_{u_{1}}, \tau_{u_{1}}\right] u_{1}$.

## 2. Preliminaries

In this section we assume that the hypotheses of Theorem 1.1 or 1.2 are satisfied. In particular, $\left(F_{1}\right)-\left(F_{4}\right)$ hold. To simplify notation, $\Omega$ will stand for $\mathbb{R}^{N}$ or for a bounded domain in $\mathbb{R}^{N}$.

Lemma 2.1. If $f(x, u) \neq 0$, then $F(x, u)<\frac{1}{2} f(x, u) u$.
Proof. Suppose $u>0$. Since $f(x, t) / t \rightarrow 0$ as $t \rightarrow 0$ and $f(x, u) / u>0$,

$$
F(x, u)=\int_{0}^{u} \frac{f(x, t)}{t} t d t<\frac{f(x, u)}{u} \int_{0}^{u} t d t=\frac{1}{2} f(x, u) u
$$

For $u<0$ the proof is similar.
The following result will be crucial for studying the structure of the set $\widehat{E}(u) \cap \mathcal{M}$.
Proposition 2.2. Let $x \in \Omega$ be fixed and let $u, s, v \in \mathbb{R}$ be such that $s \geq 0$ and $f(x, u) \neq 0$. Then:

$$
\begin{equation*}
g(s, v):=f(x, u)\left[\frac{1}{2}\left(s^{2}-1\right) u+s v\right]+F(x, u)-F(x, s u+v) \leq 0 \tag{i}
\end{equation*}
$$

for all $x$.
(ii) There exist $s_{u} \in(0,1], t_{u} \geq 1$ such that $g(s, v)=0$ if and only if $s \in\left[s_{u}, t_{u}\right]$ and $v=0$ ( $s_{u}=t_{u}$ not excluded). Moreover, for such s we have $f(x, s u)=s f(x, u)$.

Part (i) of this proposition has been shown in [7] and it extends a similar result in [11] where $\left(F_{4}^{\prime}\right)$ has been assumed (however, our $s$ corresponds to $s+1$ in $[7,11]$ ). Here we provide a different argument which will be needed in order to show part (ii).

Proof. Obviously, $g(1,0)=0$. We shall show that $g(s, v) \rightarrow-\infty$ as $s+|v| \rightarrow \infty$. Put $z=z(s):=s u+v$. Using Lemma 2.1, we obtain

$$
\begin{aligned}
g(s, v) & =f(x, u)\left[\frac{1}{2}\left(s^{2}-1\right) u+s v\right]+F(x, u)-F(x, z) \\
& <f(x, u)\left[\frac{1}{2}\left(s^{2}-1\right) u+s(z-s u)\right]+\frac{1}{2} f(x, u) u-F(x, z) \\
& =-\frac{1}{2} s^{2} f(x, u) u+s f(x, u) z-A z^{2}+\left(A z^{2}-F(x, z)\right) .
\end{aligned}
$$

Since the quadratic form (in $s$ and $z$ ) above is negative definite if $A>0$ is a constant large enough and since $A z^{2}-F(x, z)$ is bounded above according to $\left(F_{3}\right), g(s, v) \rightarrow-\infty$ as $s+|v| \rightarrow \infty$ as claimed.

It follows that $g$ has a maximum $\geq 0$ on the set $\{(s, v): s \geq 0\}$. As

$$
g(0, v)=-\frac{1}{2} f(x, u) u+F(x, u)-F(x, v)<-F(x, v) \leq 0
$$

(by Lemma 2.1), the maximum is attained at some $(s, v)$ with $s>0$. Then

$$
\begin{equation*}
g_{v}^{\prime}(s, v)=s f(x, u)-f(x, s u+v)=0 \tag{2.2}
\end{equation*}
$$

and

$$
\begin{equation*}
g_{s}^{\prime}(s, v)=(s u+v) f(x, u)-u f(x, s u+v)=0 . \tag{2.3}
\end{equation*}
$$

Using (2.2) in (2.3) we obtain $v f(x, u)=0$. Hence $v=0$ and

$$
g_{s}^{\prime}(s, 0)=s u^{2}\left(\frac{f(x, u)}{u}-\frac{f(x, s u)}{s u}\right)=0 .
$$

By $\left(F_{4}\right)$, there must exist $s_{u}, t_{u}$ such that $s_{u} \in(0,1], t_{u} \geq 1$ and $g_{s}^{\prime}(s, 0)=0$ if and only if $s \in\left[s_{u}, t_{u}\right]$. For such $s$ we have $g(s, 0)=g(1,0)=0$ and $f(x, s u)=s f(x, u)$.

Corollary 2.3. Suppose $u \in \mathcal{M}$ and let $s \geq 0, v \in E^{0} \oplus E^{-}$. Then

$$
\int_{\Omega}\left(f(x, u)\left[\frac{1}{2}\left(s^{2}-1\right) u+s v\right]+F(x, u)-F(x, s u+v)\right) d x \leq 0
$$

and there exist $0<s_{u} \leq 1 \leq t_{u}$ such that equality holds if and only if $s \in\left[s_{u}, t_{u}\right], v=0$. Moreover, for such $s$ and almost all $x \in \Omega, f(x, s u)=s f(x, u)$.

Proof. If $u \in \mathcal{M}$, then $f(x, u(x)) \neq 0$ for $x$ on a set of positive measure. According to Proposition 2.2, inequality (2.1) holds for such $x$ and there exist $s_{u(x)} \in(0,1], t_{u(x)} \geq 1$ such that the left-hand side of (2.1) is zero if and only if $s \in\left[s_{u(x)}, t_{u(x)}\right]$ and $v(x)=0$. Moreover, for such $s, f(x, s u(x))=s f(x, u(x))$. Now one takes $s_{u}:=\operatorname{ess} \sup \left\{s_{u(x)}\right.$ : $f(x, u(x)) \neq 0\}$ and $t_{u}:=\operatorname{ess} \inf \left\{t_{u(x)}: f(x, u(x)) \neq 0\right\}$.

Note that if $f(x, u(x))=0$, then $F(x, u(x))=\int_{0}^{u(x)} f(x, t) d t=0$ because $f(x, t)=0$ for $t$ between 0 and $u(x)$ according to $\left(F_{4}\right)$. Hence the integrand above is $\leq 0$ also in this case.

Proposition 2.4. (i) If $u \in E \backslash\left(E^{0} \oplus E^{-}\right)$, then $\widehat{E}(u) \cap \mathcal{M} \neq \emptyset$.
(ii) If $w \in \widehat{E}(u) \cap \mathcal{M}$, then there exist $0<s_{w} \leq 1 \leq t_{w}$ such that $\widehat{E}(u) \cap \mathcal{M}=\left[s_{w}, t_{w}\right] w$. Moreover, $\Phi(s w)=\Phi(w), \Phi^{\prime}(s w)=s \Phi^{\prime}(w)$ for all $s \in\left[s_{w}, t_{w}\right]$ and $\Phi(z)<\Phi(w)$ for all other $z \in \widehat{E}(u)$.
(iii) $\mathcal{M}$ is bounded away from $E^{0} \oplus E^{-}$, closed and $c:=\inf _{w \in \mathcal{M}} \Phi(w)>0$. Moreover, $\left.\Phi\right|_{\mathcal{M}}$ is coercive, i.e., $\Phi(u) \rightarrow \infty$ as $u \in \mathcal{M}$ and $\|u\| \rightarrow \infty$.

Note that an immediate consequence is that if $w$ is a critical point of $\Phi$, then the whole line segment $\left[s_{w}, t_{w}\right] w$ consists of critical points.

Proof. (i) The conclusion can be found in [11, Lemma 2.6 and Theorem 3.1], see also [12, Proposition 39]. The proof is by showing that $\Phi(z) \leq 0$ for $z \in \widehat{E}(u)$ and $\|z\|$ large enough, and then weak upper semicontinuity of $\left.\Phi\right|_{\widehat{E}(u)}$ implies that there exists a positive maximum.
(ii) For each $z \in \widehat{E}(u)$ we have $z=s w+v$, where $s \geq 0$ and $v=v^{0}+v^{-} \in E^{0} \oplus E^{-}$. It has been shown in the course of the proof of [11, Proposition 2.3] and [12, Proposition 39] that

$$
\begin{aligned}
\Phi(z) & -\Phi(w)=\Phi(s w+v)-\Phi(w)=-\frac{1}{2}\left\|v^{-}\right\|^{2} \\
& +\int_{\Omega}\left(f(x, w)\left[\frac{1}{2}\left(s^{2}-1\right) w+s v\right]+F(x, w)-F(x, s w+v)\right) d x
\end{aligned}
$$

(again, keep in mind that our $s$ corresponds to $s+1$ in [11, 12]). Hence according to Corollary 2.3, $\Phi(z) \leq \Phi(w)$ for all $z \in \widehat{E}(u)$ and $\Phi(z)=\Phi(w)$ if and only if $z \in\left[s_{w}, t_{w}\right] w$. That $\Phi(s w)=\Phi(w)$ for $s \in\left[s_{w}, t_{w}\right]$ is clear and since $\Phi(s w)=\max _{\widehat{E}(u)} \Phi(z)$, it is also clear
that $\widehat{E}(u) \cap \mathcal{M}=\left[s_{w}, t_{w}\right] w$ and $\Phi(z)<\Phi(w)$ for other $z$. The equality $\Phi^{\prime}(s w)=s \Phi^{\prime}(w)$ follows immediately from the fact that $f(x, s w)=s f(x, w)$.
(iii) That $c>0$ has been shown in [11, Lemma 2.4] and is an immediate consequence of the fact that $\Phi(u)=\frac{1}{2}\|u\|^{2}+o\left(\|u\|^{2}\right)$ as $u \rightarrow 0, u \in E^{+}$. Since $\left.\Phi\right|_{E^{0} \oplus E^{-}} \leq 0, \mathcal{M}$ is bounded away from $E^{0} \oplus E^{-}$and hence closed. Finally, according to Proposition 2.7 and the proof of Theorem 3.1 in [11], $\left.\Phi\right|_{\mathcal{M}}$ is coercive.

Remark 2.5. If $f$ satisfies $\left(F_{1}\right)-\left(F_{4}\right)$ and is of the form $f(x, u)=a(x) h(u)$, where $h(u) \neq 0$ for $u \neq 0$, then $s_{w}=t_{w}=1$ in Proposition 2.4, i.e. $\widehat{E}(u)$ intersects $\mathcal{M}$ at a unique point. Assuming the contrary, suppose $t_{w}>1$ and $w>0$ on a set of positive measure (other cases are treated similarly). So meas $\{x: w(x)>d\}$ is positive for some $d>0$. We claim that $h(t) / t$ is constant for $0<t<d$. Otherwise there exist $s \in\left(1, t_{w}\right], t_{0}$ and $\varepsilon>0$ such that $\varepsilon<t_{0}<d-\varepsilon$ and

$$
\frac{h(t)}{t}<\frac{h(s t)}{s t} \quad \text { for all } t \in\left(t_{0}-\varepsilon, t_{0}+\varepsilon\right)
$$

Since the sets $\left\{x: w(x)>t_{0}+\varepsilon\right\}$ and $\left\{x: w(x)<t_{0}-\varepsilon\right\}$ have positive measure, so does the set $\left\{x: w(x) \in\left(t_{0}-\varepsilon, t_{0}+\varepsilon\right)\right\}$, see [1]. But this contradicts the last statement of Corollary 2.3. Hence $h(t) / t$ is constant for $0<t<d$ and $h(t) / t \rightarrow 0$ as $t \rightarrow 0$. So $h(t)=0$ on $(0, d)$ which is impossible.

According to Proposition 2.4, for each $u \in E^{+} \backslash\{0\}$ there exist $w$ and $0<\sigma_{w} \leq \tau_{w}$ such that

$$
m(u):=\left[\sigma_{w}, \tau_{w}\right] w=\widehat{E}(u) \cap \mathcal{M} \subset E
$$

This is a multivalued map from $E^{+} \backslash\{0\}$ to $E$. However, the map $\widehat{\Psi}: E^{+} \backslash\{0\} \rightarrow \mathbb{R}$ given by

$$
\widehat{\Psi}(u):=\Phi(m(u))=\max _{z \in \widehat{E}(u)} \Phi(z)
$$

is single-valued because $\Phi$ is constant on $\widehat{E}(u) \cap \mathcal{M}$. In fact more is true:
Proposition 2.6. The map $\widehat{\Psi}$ is locally Lipschitz continuous.
Proof. If $u_{0} \in E^{+} \backslash\{0\}$, then there exist a neighbourhood $U \subset E^{+} \backslash\{0\}$ of $u_{0}$ and $R>0$ such that $\Phi(w) \leq 0$ for all $u \in U$ and $w \in \widehat{E}(u),\|w\| \geq R$. For otherwise we can find sequences $\left(u_{n}\right),\left(w_{n}\right)$ such that $u_{n} \rightarrow u_{0}, w_{n} \in \widehat{E}\left(u_{n}\right), \Phi\left(w_{n}\right)>0$ and $\left\|w_{n}\right\| \rightarrow \infty$. But $u_{0}, u_{1}, u_{2}, \ldots$ is a compact set, hence according to [11, Lemma 2.5], $\Phi(w) \leq 0$ for some $R$ and all $w \in \widehat{E}\left(u_{j}\right), j=0,1,2, \ldots,\|w\| \geq R$, which is a contradiction.

Let $U, R$ be as above and $s_{1} u_{1}+v_{1} \in m\left(u_{1}\right), s_{2} u_{2}+v_{2} \in m\left(u_{2}\right)$, where $u_{1}, u_{2} \in U$ and $v_{1}, v_{2} \in E^{0} \oplus E^{-}$. Then $\left\|m\left(u_{1}\right)\right\|,\left\|m\left(u_{2}\right)\right\| \leq R$. By the maximality property of $m(u)$ and the mean value theorem,

$$
\begin{aligned}
\widehat{\Psi}\left(u_{1}\right)-\widehat{\Psi}\left(u_{2}\right) & =\Phi\left(s_{1} u_{1}+v_{1}\right)-\Phi\left(s_{2} u_{2}+v_{2}\right) \leq \Phi\left(s_{1} u_{1}+v_{1}\right)-\Phi\left(s_{1} u_{2}+v_{1}\right) \\
& \leq s_{1} \sup _{t \in[0,1]}\left\|\Phi^{\prime}\left(s_{1}\left(t u_{1}+(1-t) u_{2}\right)+v_{1}\right)\right\|\left\|u_{1}-u_{2}\right\| \leq C\left\|u_{1}-u_{2}\right\|
\end{aligned}
$$

where the constant $C$ depends on $R$ but not on the particular choice of points in $m\left(u_{1}\right)$, $m\left(u_{2}\right)$. Similarly, $\widehat{\Psi}\left(u_{2}\right)-\widehat{\Psi}\left(u_{1}\right) \leq C\left\|u_{1}-u_{2}\right\|$ and the conclusion follows.

Remark 2.7. It has been shown in [11] that if $\left(F_{4}^{\prime}\right)$ holds instead of $\left(F_{4}\right)$, then $\widehat{\Psi} \in$ $C^{1}\left(E^{+} \backslash\{0\}, \mathbb{R}\right)$. An easy inspection of the arguments in [11] or [12] shows that if for each $u \in E^{+} \backslash\{0\}$ there exists a unique positive maximum of $\left.\Phi\right|_{\widehat{E}(u)}$, then $\widehat{\Psi}$ is still of class $C^{1}$. Hence in particular, if $f$ is as in Remark 2.5, then the conclusions of Theorems 1.1 and 1.2 hold with the same proofs as in [11].

However, under our assumptions we can in general only assert that $\widehat{\Psi}$ is locally Lipschitz continuous (because $u \mapsto m(u)$ may not be single-valued). Therefore, instead of the derivative of $\widehat{\Psi}$ we shall use Clarke's subdifferential [4]. The study of minimax methods for differential equations whose associated functional is merely locally Lipschitz continuous has been initiated by Chang in [3]. We recall some notions and facts taken from [3, 4]. They may also be found conveniently collected in Section 7.1 of [2]. The generalized directional derivative of $\widehat{\Psi}$ at $u$ in the direction $v$ is defined by

$$
\widehat{\Psi}^{\circ}(u ; v):=\limsup _{\substack{h \rightarrow 0 \\ t \downarrow 0}} \frac{\widehat{\Psi}(u+h+t v)-\widehat{\Psi}(u+h)}{t}
$$

The function $v \mapsto \widehat{\Psi}^{\circ}(u ; v)$ is convex and its subdifferential $\partial \widehat{\Psi}(u)$ is called the generalized gradient (or Clarke's subdifferential) of $\widehat{\Psi}$ at $u$, that is,

$$
\begin{equation*}
\partial \widehat{\Psi}(u):=\left\{w \in E^{+}: \widehat{\Psi}^{\circ}(u ; v) \geq\langle w, v\rangle \text { for all } v \in E^{+}\right\} \tag{2.4}
\end{equation*}
$$

In [2] $E$ is a Banach space and the generalized gradient is in the dual space $E^{*}$. Since here we work in a Hilbert space, we may assume via duality that $\partial \widehat{\Psi}(u)$ is a subset of $E$ (or more precisely, of $E^{+}$). A point $u$ is called a critical point of $\widehat{\Psi}$ if $0 \in \partial \widehat{\Psi}(u)$, i.e. $\widehat{\Psi}^{0}(u ; v) \geq 0$ for all $v \in E^{+}$, and a sequence $\left(u_{n}\right)$ is called a Palais-Smale sequence for $\widehat{\Psi}$ (PS-sequence for short) if $\widehat{\Psi}\left(u_{n}\right)$ is bounded and there exist $w_{n} \in \partial \widehat{\Psi}\left(u_{n}\right)$ such that $w_{n} \rightarrow 0$. The functional $\widehat{\Psi}$ satisfies the $P S$-condition if each PS-sequence has a convergent subsequence. Below we collect some notation which we shall need:

$$
\begin{gathered}
S^{+}:=\left\{u \in E^{+}:\|u\|=1\right\}, \quad T_{u} S^{+}:=\left\{v \in E^{+}:\langle u, v\rangle=0\right\}, \quad \Psi:=\left.\widehat{\Psi}\right|_{S^{+}}, \\
\Psi^{d}:=\left\{u \in S^{+}: \Psi(u) \leq d\right\}, \quad \Psi_{c}:=\left\{u \in S^{+}: \Psi(u) \geq c\right\}, \quad \Psi_{c}^{d}:=\Psi_{c} \cap \Psi^{d}, \\
K:=\left\{u \in S^{+}: 0 \in \partial \widehat{\Psi}(u)\right\} \quad K_{c}:=\Psi_{c}^{c} \cap K, \quad \partial \Psi(u):=\partial \widehat{\Psi}(u), \text { where } u \in S^{+} .
\end{gathered}
$$

Note that the symbol $\partial \Psi(u)$ stands for $\partial \widehat{\Psi}(u)$ when $u$ is restricted to $S^{+}$. This is in consistence with the notation $\Psi=\left.\widehat{\Psi}\right|_{S^{+}}$. As we shall see in the proof of the next proposition, $\widehat{\Psi}^{\circ}(u ; s u)=0$ for all $s \in \mathbb{R}$. Hence $\partial \Psi(u) \subset T_{u} S^{+}$.

Proposition 2.8. (i) $u \in S^{+}$is a critical point of $\widehat{\Psi}$ if and only if $m(u)$ consists of critical points of $\Phi$. The corresponding critical values coincide.
(ii) $\left(u_{n}\right) \subset S^{+}$is a PS-sequence for $\widehat{\Psi}$ if and only if there exist $w_{n} \in m\left(u_{n}\right)$ such that $\left(w_{n}\right)$ is a PS-sequence for $\Phi$.

Proof. (i) Let $u \in S^{+}$. We shall show that $\widehat{\Psi}^{\circ}(u ; v) \geq 0$ for all $v \in E^{+}$if and only if $m(u)$ consists of critical points. Note first that there exists an orthogonal decomposition $E=E(u) \oplus T_{u} S^{+}$, and by the maximizing property of $m(u), \Phi^{\prime}(w) v=0$ for all $w \in m(u)$ and $v \in E(u)$. Let $s \in \mathbb{R}$ be fixed. Since $\widehat{\Psi}(u)=\widehat{\Psi}(\sigma u)$ for all $\sigma>0$ and $\widehat{\Psi}$ is locally Lipschitz continuous,

$$
|\widehat{\Psi}(u+h+t(s u))-\widehat{\Psi}(u+h)|=|\widehat{\Psi}((1+t s) u+h)-\widehat{\Psi}((1+t s)(u+h))| \leq C t|s|\|h\|
$$

for $\|h\|$ and $t>0$ small. Hence $\widehat{\Psi}^{\circ}(u ; s u)=0$ for all $s \in \mathbb{R}$. So we only need to consider $v \in T_{u} S^{+}$.

Let $s_{u} u+z_{u}$, where $s_{u}>0$ and $z_{u} \in E^{0} \oplus E^{-}$, denote an (arbitrarily chosen) element of $m(u)$. Then, using the maximizing property of $m(u)$ and the mean value theorem,

$$
\begin{aligned}
\widehat{\Psi}(u+h+t v) & -\widehat{\Psi}(u+h)=\Phi\left(s_{u+h+t v}(u+h+t v)+z_{u+h+t v}\right)-\Phi\left(s_{u+h}(u+h)+z_{u+h}\right) \\
& \leq \Phi\left(s_{u+h+t v}(u+h+t v)+z_{u+h+t v}\right)-\Phi\left(s_{u+h+t v}(u+h)+z_{u+h+t v}\right) \\
& =t s_{u+h+t v} \Phi^{\prime}\left(s_{u+h+t v}(u+h+\theta t v)+z_{u+h+t v}\right) v
\end{aligned}
$$

for some $\theta \in(0,1)$. Dividing by $t$ and letting $h \rightarrow 0$ and $t \downarrow 0$ via subsequences we obtain

$$
\begin{equation*}
\widehat{\Psi}^{\circ}(u ; v) \leq s^{*} \Phi^{\prime}\left(s^{*} u+z^{*}\right) v \tag{2.5}
\end{equation*}
$$

where $s_{n}:=s_{u+h_{n}+t_{n} v} \rightarrow s^{*}>0$ and $z_{n}:=z_{u+h_{n}+t_{n} v} \rightharpoonup z^{*}$. This follows because $\mathcal{M}$ is bounded away from 0 and $\left.\Phi\right|_{\mathcal{M}}$ coercive, hence $s_{n}$ and $z_{n}$ must be bounded. We claim that $s^{*} u+z^{*} \in \mathcal{M}$. Indeed, taking subsequences once more, writing $z_{n}=z_{n}^{0}+z_{n}^{-} \in E^{0} \oplus E^{-}$ and using Fatou's lemma,

$$
\begin{aligned}
\widehat{\Psi}(u) & =\lim _{n \rightarrow \infty} \widehat{\Psi}\left(u+h_{n}+t_{n} v\right)=\lim _{n \rightarrow \infty} \Phi\left(s_{n}\left(u+h_{n}+t_{n} v\right)+z_{n}\right) \\
& =\lim _{n \rightarrow \infty}\left(\frac{1}{2}\left\|s_{n}\left(u+h_{n}+t_{n} v\right)\right\|^{2}-\frac{1}{2}\left\|z_{n}^{-}\right\|^{2}-\int_{\Omega} F\left(x, s_{n}\left(u+h_{n}+t_{n} v\right)+z_{n}\right) d x\right) \\
& \leq \frac{1}{2}\left\|s^{*} u\right\|^{2}-\frac{1}{2}\left\|\left(z^{*}\right)^{-}\right\|^{2}-\int_{\Omega} F\left(x, s^{*} u+z^{*}\right) d x \leq \widehat{\Psi}(u) .
\end{aligned}
$$

This implies that $\left\|z_{n}\right\| \rightarrow\left\|z^{*}\right\|$ (recall $\operatorname{dim} E^{0}<\infty$ ), hence $z_{n} \rightarrow z^{*}$ and $s_{n}\left(u+h_{n}+\right.$ $\left.t_{n} v\right)+z_{n} \rightarrow s^{*} u+z^{*}$. As $\mathcal{M}$ is closed, the claim follows. Since $\widehat{E}(u) \cap \mathcal{M}$ may be a line segment, it is not sure that $s^{*}$ and $z^{*}$ are the same for different $v$. However, if $s_{1}^{*}, s_{2}^{*}$ and $z_{1}^{*}, z_{2}^{*}$ correspond to $v_{1}$ and $v_{2}$, then by Proposition 2.4, $s_{1}^{*} u+z_{1}^{*}=\tau\left(s_{2}^{*} u+z_{2}^{*}\right)$ and $\Phi^{\prime}\left(s_{1}^{*} u+z_{1}^{*}\right) v_{2}=\tau \Phi^{\prime}\left(s_{2}^{*} u+z_{2}^{*}\right) v_{2}$ for some $\tau>0$. Taking this into account, we see from (2.5) that if $y \in \partial \Psi(u)$, then

$$
\begin{equation*}
\langle y, v\rangle \leq \widehat{\Psi}^{\circ}(u ; v) \leq \tau(v) \Phi^{\prime}\left(s^{*} u+z^{*}\right) v \tag{2.6}
\end{equation*}
$$

where $\tau$ is bounded and bounded away from 0 (by constants independent of $v$ ). It follows immediately that $u$ is a critical point of $\Psi$ if and only if $m(u)$ consists of critical points of $\Phi$.
(ii) The proof is very similar here. We take $y_{n} \in \partial \Psi\left(u_{n}\right)$ and $w_{n} \in m\left(u_{n}\right)$. Since $\left.\Phi\right|_{\mathcal{M}}$ is coercive, boundedness of $\Phi\left(m\left(u_{n}\right)\right)$ implies that $\left(m\left(u_{n}\right)\right)$ is bounded. As in (2.6), we
see that

$$
\begin{equation*}
\left\langle y_{n}, v\right\rangle \leq \widehat{\Psi}^{\circ}\left(u_{n} ; v\right) \leq \tau_{n}(v) \Phi^{\prime}\left(w_{n}\right) v, \tag{2.7}
\end{equation*}
$$

where $\tau_{n}$ is bounded and bounded away from 0 because so is $m\left(u_{n}\right)$. So the conclusion follows.

Note that if $\left(w_{n}\right) \subset\left(m\left(u_{n}\right)\right)$ is a PS-sequence for $\Phi$, then so is any sequence $\left(w_{n}^{\prime}\right) \subset$ ( $m\left(u_{n}\right)$ ).

Finally for this section we construct a pseudo-gradient vector field $H: S^{+} \backslash K \rightarrow T S^{+}$ for $\Psi$. For $u \in S^{+}$, let
$\partial^{-} \Psi(u):=\left\{p \in \partial \Psi(u):\|p\|=\min _{a \in \partial \Psi(u)}\|a\|\right\} \quad$ and $\quad \mu(u):=\inf _{a \in S^{+}}\left\{\left\|\partial^{-} \Psi(a)\right\|+\|u-a\|\right\}$.
Since $\partial \Psi(u)$ is closed and convex, $p$ as above exists and is unique, cf. [2, 3]. Hence

$$
K=\left\{u \in S^{+}: \partial^{-} \Psi(u)=0\right\} .
$$

The map $u \mapsto\left\|\partial^{-} \Psi(u)\right\|$ is lower semicontinuous [2, Proposition 7.1.1(vi)] but not continuous in general. The reason for introducing the function $\mu$ is that it regularizes $\left\|\partial^{-} \Psi(u)\right\|$. The idea comes from [5] where a similar function has been defined.

Lemma 2.9. The function $\mu$ is continuous and $u \in K$ if and only if $\mu(u)=0$.
Proof. Let $u, v, a \in S^{+}$. Then

$$
\mu(u) \leq\left\|\partial^{-} \Psi(a)\right\|+\|u-a\| \leq\left\|\partial^{-} \Psi(a)\right\|+\|v-a\|+\|u-v\|,
$$

and taking the infimum over $a$ on the right-hand side we obtain $\mu(u) \leq \mu(v)+\|u-v\|$. Reversing the roles of $u$ and $v$ we see that $|\mu(u)-\mu(v)| \leq\|u-v\|$. Hence $\mu$ is (Lipschitz) continuous.

Since $0 \leq \mu(u) \leq\left\|\partial^{-} \Psi(u)\right\|$, it is clear that $\mu(u)=0$ if $u \in K$. Suppose $\mu(u)=0$. Then there exist $a_{n}$ such that $\partial^{-} \Psi\left(a_{n}\right) \rightarrow 0$ and $a_{n} \rightarrow u$, so $u \in K$ by the lower semicontinuity of $u \mapsto\left\|\partial^{-} \Psi(u)\right\|$.

Proposition 2.10. There exists a locally Lipschitz continuous function $H: S^{+} \backslash K \rightarrow T S^{+}$ such that $\|H(u)\| \leq 1$ and $\inf \{\langle p, H(u)\rangle: p \in \partial \Psi(u)\}>\frac{1}{2} \mu(u)$ for all $u \in S^{+} \backslash K$. If $\Phi$ is even, then $H$ may be chosen to be odd.

Proof. Let $u \in S^{+} \backslash K$ and put $v_{u}:=\partial^{-} \Psi(u) /\left\|\partial^{-} \Psi(u)\right\|$. Consider the map

$$
\chi: w \mapsto \inf _{p \in \partial \Psi(w)}\left\langle p, v_{u}-\left\langle v_{u}, w\right\rangle w\right\rangle-\frac{1}{2} \mu(w), \quad w \in S^{+} \backslash K
$$

(note that $v_{u}-\left\langle v_{u}, w\right\rangle w \in T_{w}\left(S^{+}\right)$). Since $\partial \Psi(u)$ is convex, $\inf _{p \in \partial \Psi(u)}\left\langle p, v_{u}\right\rangle \geq\left\|\partial^{-} \Psi(u)\right\| \geq$ $\mu(u)$ and therefore $\chi(u) \geq \frac{1}{2} \mu(u)>0$. Moreover, since

$$
\inf _{p \in \partial \Psi(w)}\left\langle p, v_{u}-\left\langle v_{u}, w\right\rangle w\right\rangle=-\sup _{p \in \partial \Psi(w)}\left\langle p,\left\langle v_{u}, w\right\rangle w-v_{u}\right\rangle=-\widehat{\Psi}^{0}\left(w ;\left\langle v_{u}, w\right\rangle w-v_{u}\right)
$$

(see Proposition 7.1.1(vii) and property (c) on p. 168 in [2]) and $\widehat{\Psi}^{0}$ is upper semicontinuous in both arguments [2, Proposition 7.1.1(vii)], $\chi$ is lower semicontinuous. Hence there exists a neighbourhood $U_{u}$ of $u$ such that $\chi(w)>0$ for all $w \in U$.

The remaining part of the proof is standard. Take a locally finite open refinement $\left(U_{u_{i}}\right)_{i \in I}$ of the open cover $\left(U_{u}\right)_{u \in S^{+} \backslash K}$ and a subordinated locally Lipschitz continuous partition of unity $\left\{\lambda_{i}\right\}_{i \in I}$. Define

$$
H(u):=\sum_{i \in I} \lambda_{i}(u) v_{u_{i}}, \quad u \in S^{+} \backslash K .
$$

It is easy to see that $H$ satisfies the required conclusions.
If $\Phi$ is even, then so is $\Psi$ and we may replace $H(u)$ with $\frac{1}{2}(H(u)-H(-u))$.

## 3. Proofs of Theorems 1.1 and 1.2

Since the arguments are very similar to those appearing in [11, 12], we shall describe them rather briefly and concentrate on pointing out the main differences.

We start with Theorem 1.1. First we want to show that there exists a minimizer for $\Psi$ on $S^{+}$. It follows from the results of Section 2 that

$$
c:=\inf _{w \in \mathcal{M}} \Phi(w)=\inf _{u \in S^{+}} \Psi(u)>0
$$

According to Ekeland's variational principle [6], there exists a sequence $\left(u_{n}\right) \subset S^{+}$such that $\Psi\left(u_{n}\right) \rightarrow c$ and

$$
\begin{equation*}
\Psi(w) \geq \Psi\left(u_{n}\right)-\frac{1}{n}\left\|w-u_{n}\right\| \quad \text { for all } w \in S^{+} \tag{3.1}
\end{equation*}
$$

For a given $v \in T_{u_{n}} S^{+}$, let $z_{n}(t):=\left(u_{n}+t v\right) /\left\|u_{n}+t v\right\|$. Since $\left\|u_{n}+t v\right\|-1=O\left(t^{2}\right)$ as $t \rightarrow 0$ and $\widehat{\Psi}\left(u_{n}+t v\right)=\Psi\left(z_{n}(t)\right)$, it follows from (3.1) that

$$
\widehat{\Psi}^{\circ}\left(u_{n} ; v\right) \geq \limsup _{t \downarrow 0} \frac{\widehat{\Psi}\left(u_{n}+t v\right)-\widehat{\Psi}\left(u_{n}\right)}{t}=\underset{t \downarrow 0}{\lim \sup } \frac{\Psi\left(z_{n}(t)\right)-\Psi\left(u_{n}\right)}{t} \geq-\frac{1}{n}\|v\| .
$$

Since $m\left(u_{n}\right)$ is bounded by coercivity of $\left.\Phi\right|_{\mathcal{M}}$, the second inequality in (2.7) implies that

$$
-\frac{1}{n}\|v\| \leq \widehat{\Psi}^{\circ}\left(u_{n} ; v\right) \leq \tau_{n}(v) \Phi^{\prime}\left(w_{n}\right) v
$$

where $w_{n} \in m\left(u_{n}\right) \subset \mathcal{M}$ and $\tau_{n}$ is bounded and bounded away from 0 . So recalling $\Phi^{\prime}\left(w_{n}\right) v=0$ for all $v \in E\left(w_{n}\right)$, it follows that $\left(w_{n}\right)$ is a bounded PS-sequence for $\Phi$. Now we may proceed exactly as in the proof of Theorem 1.1 in [11], pp. 3811-3812 (or in the proof of Theorem 40 in [12]). More precisely, one shows invoking Lions' lemma [15, Lemma 1.21] in a rather standard way that there exists a sequence $\left(y_{n}\right) \subset \mathbb{R}^{N}$ such that

$$
\int_{\left|x-y_{n}\right|<1} w_{n}^{2} d x \geq \varepsilon \text { for } n \text { large enough and some } \varepsilon>0
$$

and since $\Phi$ and $\mathcal{M}$ are invariant by translations $u(\cdot) \mapsto u(\cdot-k), k \in \mathbb{Z}^{N}$, we may assume $\left(y_{n}\right)$ is bounded. So passing to a subsequence, $w_{n} \rightharpoonup w \neq 0$. This $w$ is a solution and an additional argument shows it is a ground state, see [11] or [12] for more details.

Suppose now $f$ is odd in $u$ and note that $\Psi$ is even and invariant by translations by elements of $\mathbb{Z}^{N}$. To prove that there exist infinitely many geometrically distinct solutions we assume the contrary. Since to each $\left[\sigma_{w}, \tau_{w}\right] w \subset \mathcal{M}$ there corresponds a unique point $u \in S^{+}, K$ consists of finitely many orbits $\mathcal{O}(u):=\left\{u(\cdot-k): u \in K, k \in \mathbb{Z}^{N}\right\}$. We choose a subset $\mathcal{F} \subset K$ such that $\mathcal{F}=-\mathcal{F}$ and each orbit has a unique representative in $\mathcal{F}$. Now an easy inspection shows that Lemmas 2.11 and 2.13 in [11] hold, i.e. the mapping $\check{m}: u \mapsto u^{+} /\left\|u^{+}\right\|$from $\mathcal{M}$ to $S^{+}$is Lipschitz continuous and $\kappa:=\inf \{v-w \|: v, v \in$ $K, v \neq w\}>0$.

Proposition 3.1 (Lemma 2.14 in [11]). Let $d \geq c$. If $\left(v_{n}^{1}\right),\left(v_{n}^{2}\right) \subset \Psi^{d}$ are two PS-sequences for $\Psi$, then either $\left\|v_{n}^{1}-v_{n}^{2}\right\| \rightarrow 0$ as $n \rightarrow \infty$ or $\lim \sup _{n \rightarrow \infty}\left\|v_{n}^{1}-v_{n}^{2}\right\| \geq \rho(d)>0$, where $\rho$ depends on d but not on the particular choice of PS-sequences in $\Psi^{d}$.

The argument is exactly the same as in [11], taking into account that by Proposition 2.8, to $\left(v_{n}^{j}\right) \subset \Psi^{d}$ there correspond PS-sequences $\left(u_{n}^{j}\right)$ with $u_{n}^{j} \in m\left(v_{n}^{j}\right), j=1,2$. Once $u_{n}^{j}$ have been chosen, one follows the lines of [11].

Let $H$ be the vector field constructed in Proposition 2.10 and consider the flow given by

$$
\frac{d}{d t} \eta(t, w)=-H(\eta(t, w)), \quad \eta(0, w)=w
$$

defined on the set

$$
\mathcal{G}:=\left\{(t, w): w \in S^{+} \backslash K, T^{-}(w)<t<T^{+}(w)\right\}
$$

where $\left(T^{-}(w), T^{+}(w)\right)$ is the maximal existence time for the trajectory passing through $w$ at $t=0$.

Proposition 3.2 (cf. Lemma 2.15 in [11]). For each $w \in S^{+} \backslash K$ the $\operatorname{limit}^{\lim }{ }_{t \rightarrow T^{+}(w)} \eta(t, w)$ exists and is a critical point of $\Psi$.

Proof. We adapt the argument in [11].
If $T^{+}(w)<\infty$, then for $0 \leq s<t<T^{+}(w)$ we have

$$
\|\eta(t, w)-\eta(s, w)\| \leq \int_{s}^{t}\|H(\eta(\tau, w))\| d \tau \leq t-s
$$

hence $\lim _{t \rightarrow T^{+}(w)} \eta(t, w)$ exists and must be a critical point (or the flow can be continued for $\left.t>T^{+}(w)\right)$.

Let $T^{+}(w)=\infty$. It suffices to show that for each $\varepsilon>0$ there exists $t_{\varepsilon}>0$ such that $\left\|\eta\left(t_{\varepsilon}, u\right)-\eta(t, u)\right\|<\varepsilon$ for all $t \geq t_{\varepsilon}$. Assuming the contrary, we find $\varepsilon \in(0, \rho(d) / 2)$ and $t_{n} \rightarrow \infty$ such that $\left\|\eta\left(t_{n}, w\right)-\eta\left(t_{n+1}, w\right)\right\|=\varepsilon$ for all $n$. Choose the smallest $t_{n}^{1} \in\left(t_{n}, t_{n+1}\right)$ such that $\left\|\eta\left(t_{n}, w\right)-\eta\left(t_{n}^{1}, w\right)\right\|=\varepsilon / 3$. Recall from Lemma 2.9 that $\mu$ is continuous and set
$\kappa_{n}:=\min _{s \in\left[t_{n}, t_{n}^{1}\right]} \mu(\eta(s, w))$. Then, using Proposition 2.10 and [2, Proposition 7.1.1(viii)],

$$
\begin{aligned}
\frac{\varepsilon}{3} & =\left\|\eta\left(t_{n}^{1}, w\right)-\eta\left(t_{n}, w\right)\right\| \leq \int_{t_{n}}^{t_{n}^{1}}\|H(\eta(s, w))\| d s \leq t_{n}^{1}-t_{n} \\
& \leq \frac{2}{\kappa_{n}} \int_{t_{n}}^{t_{n}^{1}} \inf _{p \in \partial \Psi(\eta(s, u))}\langle p, H(\eta(s, w))\rangle d s=-\frac{2}{\kappa_{n}} \int_{t_{n}}^{t_{n}^{1}} \sup _{p \in \partial \Psi(\eta(s, u))}\langle p,-H(\eta(s, w))\rangle d s \\
& \leq-\frac{2}{\kappa_{n}} \int_{t_{n}}^{t_{n}^{1}} \frac{d}{d s} \Psi(\eta(s, w)) d s=\frac{2}{\kappa_{n}}\left(\Psi\left(\eta\left(t_{n}, w\right)\right)-\Psi\left(\eta\left(t_{n}^{1}, w\right)\right)\right) .
\end{aligned}
$$

Since $\Psi$ is bounded below, $\Psi\left(\eta\left(t_{n}, w\right)\right)-\Psi\left(\eta\left(t_{n}^{1}, w\right)\right) \rightarrow 0$, hence $\kappa_{n} \rightarrow 0$ and we may find $s_{n}^{1} \in\left[t_{n}, t_{n}^{1}\right]$ such that if $z_{n}^{1}:=\eta\left(s_{n}^{1}, w\right)$, then $\mu\left(z_{n}^{1}\right) \rightarrow 0$. By the definition of $\mu$ there exist $w_{n}^{1}$ such that $w_{n}^{1}-z_{n}^{1} \rightarrow 0$ and $\partial^{-} \Psi\left(w_{n}^{1}\right) \rightarrow 0$. So $\lim \sup _{n \rightarrow \infty}\left\|w_{n}^{1}-\eta\left(t_{n}, w\right)\right\| \leq \varepsilon / 3$. Similarly, there exists a largest $t_{n}^{2} \in\left(t_{n}^{1}, t_{n+1}\right)$ with $\left\|\eta\left(t_{n+1}, w\right)-\eta\left(t_{n}^{2}, w\right)\right\|=\varepsilon / 3$ and we find $w_{n}^{2}$ with $\partial^{-} \Psi\left(w_{n}^{2}\right) \rightarrow 0$ and $\lim \sup _{n \rightarrow \infty}\left\|w_{n}^{2}-\eta\left(t_{n+1}, w\right)\right\| \leq \varepsilon / 3$. It follows that $\varepsilon / 3 \leq \lim \sup _{n \rightarrow \infty}\left\|w_{n}^{1}-w_{n}^{2}\right\| \leq 2 \varepsilon<\rho(d)$, a contradiction to Proposition 3.1.

Proposition 3.3 (cf. Lemma 2.16 in [11]). Let $d \geq c$. Then for each $\delta>0$ there exists $\varepsilon>0$ such that $\Psi_{d-\varepsilon}^{d+\varepsilon} \cap K=K_{d}$ and $\lim _{t \rightarrow T^{+}(w)} \Psi(\eta(t, w))<d-\varepsilon$ for all $w \in \Psi^{d+\varepsilon} \backslash U_{\delta}\left(K_{d}\right)$, where $U_{\delta}\left(K_{d}\right)$ is the open $\delta$-neighbourhood of $K_{d}$.

The proof requires changes which, in view of the arguments of Proposition 3.2, are rather obvious (in particular, $\nabla \Psi$ in the definition of $\tau$ in [11] should be replaced by $\mu$ ).

With all these prerequisites, existence of infinitely many solutions is obtained by repeating the arguments on pp. 3817-3818 in [11]. Let

$$
c_{k}:=\inf \left\{d \in \mathbb{R}: \gamma\left(\Psi^{d}\right) \geq k\right\}, \quad k=1,2, \ldots,
$$

where $\gamma$ denotes Krasnoselskii's genus [10]. Using the flow $\eta$ and Proposition 3.3 one shows $K_{c_{k}} \neq 0$ and $c_{k}<c_{k+1}$ for all $k$. This contradicts our assumption that there are finitely many geometrically distinct solutions.

Now we turn our attention to Theorem 1.2. Here there is no $\mathbb{Z}^{N}$-symmetry but instead there is a compact embedding $H_{0}^{1}(\Omega) \hookrightarrow L^{q}(\Omega)$ for $q \in\left[1,2^{*}\right)$. Using this, one sees as in the proof of Theorem 3.1 in [11] or Theorem 37 in [12] that the ground state exists. The minimizing PS-sequence is extracted by using Ekeland's variational principle in the same way as at the beginning of this section. To obtain infinitely many solutions for odd $f$ one first shows as in [11, Theorem 3.2] (or in [12, Section 4.2]) that $\Psi$ satisfies the PScondition. Now a standard minimax argument as in [11, Theorem 3.2] can be employed. Note that with the aid of the vector field $H$ and suitable cutoff functions one can construct a deformation in the usual way as e.g. in [10] (see also [3]). We leave the details to the reader.

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